

# Fast Diffusion Equation via Diffusive Scaling of Generalized Carleman Kinetic Equation

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# Overview

- 1 Introduction
  - Model and History
  - Notations
  - Formal Computation
- 2 Flux Estimate via Local Entropy
- 3 Diffusive Limit with General Initial Data
- 4 Comparison Principle
- 5 Main Theorem

# Introduction-Model

Discrete velocity kinetic equation of a fictitious gas proposed by [Carleman '33]. For  $|\alpha| \leq 1$ ,

$$\begin{cases} \partial_t u_i^\varepsilon + \frac{1}{\varepsilon} \nabla_i u_i^\varepsilon = \frac{1}{2n\varepsilon^2} \sum_{j=1}^{2n} (u_j^\varepsilon + u_i^\varepsilon)^\alpha (u_j^\varepsilon - u_i^\varepsilon) & \text{in } Q \\ \partial_t u_{i+n}^\varepsilon - \frac{1}{\varepsilon} \nabla_i u_{i+n}^\varepsilon = \frac{1}{2n\varepsilon^2} \sum_{j=1}^{2n} (u_j^\varepsilon + u_{i+n}^\varepsilon)^\alpha (u_j^\varepsilon - u_{i+n}^\varepsilon) & \text{in } Q \end{cases} \quad (1.1)$$

with initial condition  $u_i^\varepsilon = g_i \geq 0$  at  $t = 0$  ( $i = 1, \dots, 2n$ ).

# Introduction-Model

$$\begin{cases} \partial_t u_i^\varepsilon + \frac{1}{\varepsilon} \nabla_i u_i^\varepsilon = \frac{1}{2n\varepsilon^2} \sum_{j=1}^{2n} (u_j^\varepsilon + u_i^\varepsilon)^\alpha (u_j^\varepsilon - u_i^\varepsilon) & \text{in } Q \\ \partial_t u_{i+n}^\varepsilon - \frac{1}{\varepsilon} \nabla_i u_{i+n}^\varepsilon = \frac{1}{2n\varepsilon^2} \sum_{j=1}^{2n} (u_j^\varepsilon + u_{i+n}^\varepsilon)^\alpha (u_j^\varepsilon - u_{i+n}^\varepsilon) & \text{in } Q \end{cases}$$

- For  $n = 1$  and  $\alpha = 1$ , proposed by Carleman
- When  $\varepsilon = 1$ ,  $u_i$  and  $u_{i+n}$  represents density of particles moving in  $x_i$  and  $-x_i$  direction with speed 1, respectively
- $(u_i + u_j)^\alpha (u_j - u_i)$  describes the interaction between two group of particles,  $i$  and  $j$  at a position  $x$

# Introduction-Model

- Eq (1.1) comes from  $(x, t) \rightarrow (\varepsilon x, \varepsilon^2 t)$ , diffusive limit
- Physically,  $\varepsilon > 0$  represents Knudsen number and Mach number
- In  $\varepsilon \rightarrow 0$  limit, this mesoscopic kinetic equation becomes a certain continuum equation

- GOAL

Prove convergence of  $\rho^\varepsilon := \sum_{i=1}^{2n} u_i^\varepsilon$  to  $\rho$  in rigorous way and show the limit  $\rho$  satisfies

$$\partial_t \rho - \nabla \cdot \left( \frac{1}{n^{1-\alpha} \cdot \rho^\alpha} \nabla \rho \right) = 0 \text{ with } \rho(0) = \sum_{i=1}^{2n} g_i \quad (1.2)$$

- Fast Diffusion Eq if  $\alpha > 0$ , and Porous Medium Eq if  $\alpha < 0$

## Regarding Diffusive Scaling Limit

- For  $n = 1$  and  $\alpha = 1$ , [Kurtz '73] and [McKean '75]
- For  $n = 1$  and  $\alpha < 1$ , [Pulvirenti&Toscani '96], [Toscani&Lions '97]
- For  $n \geq 1$  and  $\alpha < 1$ , [Toscani&Lions '97] slightly different model
- For  $n = 1$  and  $1 < \alpha < \frac{4}{3}$ , [Salvarani&Toscani '09]
- For  $n = 1$  and  $|\alpha| \leq 1$ , [Salvarani&Vazquez '05] with general  $L^1 + L^\infty$  initial data  $g_i$  without weight or flux assumption

Previous results use entropy method

- Uniform estimates (lower bound) on functionals such as  $\int \sum_{i=1}^{2n} u_i^\varepsilon \log u_i^\varepsilon dx$  or  $\int \sum_{i=1}^{2n} \phi(u_i^\varepsilon) dx$  allow us to pass the limit
- To obtain such an estimate, uniform estimate of a kind of weighted  $L^1$  norm is required
- This is undesirable if target fast diffusion equation is not mass preserving
- In higher dimension, some fast diffusion is no longer mass preserving and a new method is needed.



# Introduction-Notations

These will be some notations that appear frequently.

- $Q = \mathbb{R}^n \times (0, \infty)$ ,  $Q_T = \mathbb{R}^n \times (0, T)$ ,  $T > 0$ ,  $Q_{R,T} = B_R(0) \times (0, T)$

- $\rho^\varepsilon = \sum_{i=1}^{2n} u_i^\varepsilon$ ,  $\rho_i^\varepsilon = u_i^\varepsilon + u_{i+n}^\varepsilon$ ,  $\rho_{i,j}^\varepsilon = u_i^\varepsilon + u_j^\varepsilon$

- $J^\varepsilon = (J_1^\varepsilon, J_2^\varepsilon, \dots, J_n^\varepsilon)$ ,  $J_i^\varepsilon = \frac{u_i^\varepsilon - u_{i+n}^\varepsilon}{\varepsilon}$ ,  $J_{i,j}^\varepsilon = \frac{u_i^\varepsilon - u_j^\varepsilon}{\varepsilon}$

$J^\varepsilon$  will behave as  $-C \frac{\nabla \rho}{\rho^\alpha}$  in  $\varepsilon \rightarrow 0$  in limit

# Introduction-Formal Computation in 1D

In 1D case, our equation could be written

$$\begin{cases} \partial_t u_1^\varepsilon + \frac{1}{\varepsilon} \nabla_x u_1^\varepsilon = \frac{1}{2\varepsilon^2} (u_1^\varepsilon + u_2^\varepsilon)^\alpha (u_2^\varepsilon - u_1^\varepsilon) \\ \partial_t u_2^\varepsilon - \frac{1}{\varepsilon} \nabla_x u_2^\varepsilon = \frac{1}{2\varepsilon^2} (u_1^\varepsilon + u_2^\varepsilon)^\alpha (u_1^\varepsilon - u_2^\varepsilon) \end{cases} \quad (1.3)$$

By adding and subtracting equations above for each  $i$ ,

$$\begin{cases} \partial_t \rho^\varepsilon + \nabla_x J^\varepsilon = 0 \\ \varepsilon^2 \partial_t J^\varepsilon + \nabla_x \rho^\varepsilon = -(\rho^\varepsilon)^\alpha J^\varepsilon. \end{cases} \quad (1.4)$$

If  $\rho^\varepsilon, J^\varepsilon$  are  $\varepsilon$ -uniformly  $L^2_{loc}$  bounded in space-time  $(x, t)$ , we obtain  $\rho^\varepsilon \rightharpoonup \rho, J^\varepsilon \rightharpoonup J$  in  $L^2_{loc}$ .

$$\begin{cases} \partial_t \rho + \nabla_x J = 0 \\ \nabla_x \rho = -(\rho)^\alpha J \end{cases} \quad (1.5)$$

$\therefore$  if  $\rho$  is nonzero in certain sense,  $\partial_t \rho = \nabla_x \left( \frac{\nabla_x \rho}{\rho^\alpha} \right)$ .

In fact, Divergence-Curl Lemma says convergence of  $\rho^\varepsilon$  is strong in  $L^2_{loc}$ , which reminds us compactness theorem in Sobolev spaces.

## $L^2$ bounds for $\rho^\varepsilon$ and $J^\varepsilon$

A bound for  $\rho^\varepsilon$  is easy. One can show  $\sum_i \|u_i^\varepsilon\|_p^p$  is decreasing in time, or use comparison principle. To control  $J^\varepsilon$ , we use entropy formula.

Multiply  $(u_i \log u_i)'$  on the equation and integrate in  $x$  we get,

$$\begin{aligned}\partial_t \int u_1^\varepsilon \log u_1^\varepsilon + u_2^\varepsilon \log u_2^\varepsilon dx &= - \int \frac{1}{\varepsilon^2} (\rho^\varepsilon)^\alpha (u_1^\varepsilon - u_2^\varepsilon) (\log u_1^\varepsilon - \log u_2^\varepsilon) dx \\ &\leq - \int (J^\varepsilon)^2 (\rho^\varepsilon)^{\alpha-1} dx \\ &\leq - \int (J^\varepsilon)^2 \|\rho^\varepsilon\|_\infty^{\alpha-1} dx.\end{aligned}$$

This implies...

- $\int u_1^\varepsilon \log u_1^\varepsilon + u_2^\varepsilon \log u_2^\varepsilon dx$  is decreasing in time
- For  $t_1 < t_2$ , if an upper bound of this entropy (above) at  $t_1$  and a lower bound at  $t_2$  are given, we obtain an uniform  $L^2([t_1, t_2] \times \mathbb{R})$  bound of  $J^\varepsilon$
- Same problem happens in higher dimension case
- [Salvarani&Vazquez '05] use local(in space) entropy which gives motivation of our work

# Flux Estimate via Local Entropy

## Proposition

Suppose  $u_i^\varepsilon$   $i=1, \dots, 2n$ , are uniformly bounded above and below, say  $0 < \frac{1}{M} \leq u_i^\varepsilon \leq M < \infty$  on  $K \times [0, T]$ . Then for all  $\phi \in C_0^\infty(\mathbb{R}^n)$  with  $\text{supp}(\phi) \subset K$ ,

$$\int_0^T \int \sum_{i,j} (u_i^\varepsilon + u_j^\varepsilon)^\alpha J_{i,j}^{\varepsilon 2} \phi^2 dx dt < C$$

for some  $C = C(\phi, M, \alpha, T, n) > 0$ , but independent of  $\varepsilon$ .

## Proof.

By multiplying  $\frac{1}{u_i^\varepsilon} \phi^2$  both sides of (1.1) and integrating w.r.t.  $x$ , we obtain

$$\begin{aligned} \partial_t \int \sum_{i=1}^{2n} \log(u_i^\varepsilon) \phi^2 dx + \frac{1}{\varepsilon} \int \sum_{i=1}^n \nabla_i \log(u_i^\varepsilon) - \nabla_i \log(u_{i+n}^\varepsilon) \phi^2 dx \\ = \frac{1}{4n\varepsilon^2} \int \sum_{i \neq j, 1 \leq i, j \leq 2n} (u_i^\varepsilon + u_j^\varepsilon)^\alpha (u_j^\varepsilon - u_i^\varepsilon) \left( \frac{1}{u_i^\varepsilon} - \frac{1}{u_j^\varepsilon} \right) \phi^2 dx. \end{aligned} \quad (2.1)$$

and integration by parts gives us

$$\begin{aligned} \partial_t \int \sum_{i=1}^{2n} \log(u_i^\varepsilon) \phi^2 dx = \frac{2}{\varepsilon} \int \sum_{i=1}^n (\log(u_i^\varepsilon) - \log(u_{i+n}^\varepsilon)) \phi \nabla_i \phi dx \\ - \frac{1}{4n} \int \sum_{i \neq j, 1 \leq i, j \leq 2n} \frac{(u_i^\varepsilon + u_j^\varepsilon)^\alpha (J_{j,i}^\varepsilon)^2}{(u_i^\varepsilon u_j^\varepsilon)} \phi^2 dx. \end{aligned} \quad (2.2)$$

...

# Diffusive Limit in Easy Cases

## Theorem

Suppose initial data  $\{g_i\}$  has uniform lower and upper bounds  $0 < N \leq g_i \leq M$  for  $i = 1, \dots, 2n$ , then the unique weak solution of (1.1)  $u_i^\varepsilon$  converge to the same limit  $\frac{\rho}{2n}$  in  $L^2_{loc}(Q)$  as  $\varepsilon \rightarrow 0$ . Here  $\rho$  is the unique weak solution of the Cauchy problem

$$\partial_t \rho - \nabla \cdot \left( \frac{1}{n^{1-\alpha} \cdot \rho^\alpha} \nabla \rho \right) = 0$$

with initial data  $\sum_{i=1}^n g_i$ .



# Diffusive Limit with General Initial Data

- Theorem requires uniform initial bounds  $N \leq g_i \leq M$
- Local uniform bounds  $N \leq u_i^\varepsilon \leq M$  on  $K \subset\subset Q$  on each  $K$  is enough
- IDEA. Assume  $\Psi(x) \leq g_i \leq M$  and prove  $u_i$  keep this shape of lower bound shape by making explicit barrier having initial shape like  $\Psi$
- To construct explicit (sub or super) solution, we find  $u_i^\varepsilon$  of the form  $u_i^\varepsilon = \frac{1}{2n}(\rho + A_i\varepsilon + B_i\varepsilon^2)$  and find  $A_i$  and  $B_i$  using formal asymptotic expansion.

## Definition

$(u_i)_{i=1}^{2n} \in C([0, T], L_{loc}^\infty(\mathbb{R}^n)) \cap L_{loc}^1(Q_T)$  is a weak subsolution of (1.1) if

$$\begin{cases} \partial_t u_i^\varepsilon + \frac{1}{\varepsilon} \nabla_i u_i^\varepsilon - \frac{1}{2n\varepsilon^2} \sum_{j=1}^{2n} (u_j^\varepsilon + u_i^\varepsilon)^\alpha (u_j^\varepsilon - u_i^\varepsilon) \leq 0 \\ \partial_t u_{i+n}^\varepsilon - \frac{1}{\varepsilon} \nabla_i u_{i+n}^\varepsilon - \frac{1}{2n\varepsilon^2} \sum_{j=1}^{2n} (u_j^\varepsilon + u_{i+n}^\varepsilon)^\alpha (u_j^\varepsilon - u_{i+n}^\varepsilon) \leq 0 \end{cases} \quad (4.1)$$

in distribution sense.

Similarly we may define supersolution of (1.1).

# Comparison Principle(2)

## Lemma ( $L^1$ -contraction)

If  $u_i$  and  $v_i$  are subsolution and supersolution of (1.1), respectively, then for  $t_2 > t_1 \geq 0$

$$\left[ \int_{B(0,R)} \sum_{i=1}^{2n} (u_i - v_i)^+ dx \right] (t_2) \leq \left[ \int_{B(0, R + \frac{t_2 - t_1}{\varepsilon})} \sum_{i=1}^{2n} (u_i - v_i)^+ dx \right] (t_1) \quad (4.2)$$

$$\left[ \int_{\mathbb{R}^n} \sum_{i=1}^{2n} (u_i - v_i)^+ dx \right] (t_2) \leq \left[ \int_{\mathbb{R}^n} \sum_{i=1}^{2n} (u_i - v_i)^+ dx \right] (t_1). \quad (4.3)$$

- $u_i \leq v_i$  on  $B(0, R + \frac{t_2 - t_1}{\varepsilon})$  at  $t = t_1$  for all  $i = 1, \dots, 2n$  implies that  $u_i \leq v_i$  on  $B(0, R)$   $t \in (t_1, t_2)$  for all  $i = 1, \dots, 2n$ .
- $u_i \leq v_i$  on  $\mathbb{R}^n$  at  $t = t_1$  for all  $i = 1, \dots, 2n$  implies that  $u_i \leq v_i$  on  $\mathbb{R}^n$   $t \in (t_1, t_2)$  for all  $i = 1, \dots, 2n$ .

After a long computation with symmetry assumption

$$A_i = - \left(\frac{n}{\rho}\right)^\alpha D_{v_i} \rho = \begin{cases} - \left(\frac{n}{\rho}\right)^\alpha \partial_i \rho, & 1 \leq i \leq n \\ \left(\frac{n}{\rho}\right)^\alpha \partial_{i-n} \rho, & n+1 \leq i \leq 2n \end{cases} \quad \text{and}$$

$$B_i = B_{i+n} = \left(\frac{n}{\rho}\right)^\alpha \partial_i \left( \left(\frac{n}{\rho}\right)^\alpha \partial_i \rho \right) - \frac{\alpha}{2\rho} \left( \left(\frac{n}{\rho}\right)^\alpha \partial_i \rho \right)^2 \quad 1 \leq i \leq n. \quad (4.4)$$

For our purpose,  $A_i \varepsilon$  and  $B_i \varepsilon^2$  need to be comparable to  $\rho$ , but usually  $|B_i \varepsilon^2| \& |A_i \varepsilon| \gg \rho$  as  $x \rightarrow \infty$ . Roughly, our "key observation" is that we need  $|B_i \varepsilon^2| \& |A_i \varepsilon| < c\rho$  only on the region  $|x| = O(\frac{1}{\varepsilon})$  since we need a local bound and equation has finite speed ( $= \frac{1}{\varepsilon}$ ) of propagation.

- This requires profile of  $\rho$  should not decay so fast as  $|x| \rightarrow \infty$
- Simple calculation says  $|x|^{-\frac{2}{\alpha}}$  optimal rate in this regards
- Tail profile of Barrenblat solution of equation  $\partial_t \rho = \Delta(\rho^{1-\alpha})$

For  $n \geq 3$ ,  $0 \leq 1 - \alpha < \frac{n-2}{n}$ , and  $R > 0$ , define

$$\Psi_{\alpha,n,R,T}(x, t) = C_{\alpha} \left( \frac{T - ct}{|x|^2 + R^2} \right)^{\frac{1}{\alpha}}, \quad c - 1 > \frac{2}{\alpha} \frac{1}{n - \frac{2}{\alpha}} \quad (4.5)$$

Then we may check

$$\begin{aligned} |A_i| &\leq c_0 \frac{\Psi}{T - ct} (|x|^2 + R^2)^{\frac{1}{2}}, & |\partial_t A_i| &\leq c_0 \frac{\Psi}{(T - ct)^2} (|x|^2 + R^2)^{\frac{1}{2}} \\ |B_i| &\leq c_0 \frac{\Psi}{(T - ct)^2} (|x|^2 + R^2), & |D_{v_i} B_i| &\leq c_0 \frac{\Psi}{(T - ct)^2} (|x|^2 + R^2)^{\frac{1}{2}} \\ |\partial_t B_i| &\leq c_0 \frac{\Psi}{(T - ct)^3} (|x|^2 + R^2), & |\Psi^{\alpha} A_i| &\leq c_0 \frac{\Psi}{(|x|^2 + R^2)^{\frac{1}{2}}} \\ |\Psi^{\alpha} B_i| &\leq c_0 \frac{\Psi}{(T - ct)} & & \text{and so on.} \end{aligned} \quad (4.6)$$

## Proposition

In above setting, if  $\bar{u}_i$  are defined explicitly using  $\Psi$  in place of  $\rho$ , there exists  $c_1 = c_1(c, n, \alpha) > 0$  s.t.

$$\frac{1}{4n}\Psi \leq \bar{u}_i^\varepsilon \leq \frac{3}{4n}\Psi \quad \text{for all } i = 1, \dots, 2n$$

on  $(|x|^2 + R^2)^{\frac{1}{2}} < \frac{c_1}{\varepsilon}(T - tc)$  and  $t \in [0, \frac{T}{c})$ .

## Proposition

For  $\Psi(x, t)$ , there is  $c_2 = c_2(c, n, \alpha) > 0$  such that  $\bar{u}_i$  is subsolution of (1.1) on  $(|x|^2 + R^2)^{\frac{1}{2}} < \frac{c_2}{\varepsilon}(T - tc)$  and  $t \in [0, \frac{T}{c})$ .

## Proposition

For  $n \geq 3$ ,  $0 \leq 1 - \alpha < \frac{n-2}{n}$ , suppose mild solution,  $u_i^\varepsilon$ , of (1.1) has initial data  $g_i \in L^1_{loc}(\mathbb{R}^n)$  with lower bound  $g_i \geq \frac{3}{4n} \Psi_{R,T}(0)$  for some  $R$  and  $T > 0$ . Then there exists universal constants  $C = C(n, \alpha) > 0$  such that for every space-time compact set  $K \subset \mathbb{R}^n \times [0, CT)$ , there exist  $\varepsilon_K > 0$  satisfying  $u_i^\varepsilon \geq \frac{1}{4n} \Psi_{R,T}$  on  $K$  for  $0 < \varepsilon < \varepsilon_K$ .



## Definition (Admissible Initial data)

$\mathcal{X}_{n,\alpha}$  is a collection of nonnegative  $g \in L^1_{loc}(\mathbb{R}^n)$  such that there exists nonnegative  $f \in L^1_{loc}(\mathbb{R}^n)$  with  $f - g \in L^1(\mathbb{R}^n)$  and  $f$  satisfies

$$\liminf_{|x| \rightarrow \infty} |x|^{\frac{2}{\alpha}} f > 0 \quad \text{if } \frac{2}{n} \leq \alpha \leq 1 \quad (5.1)$$

and

$$\begin{aligned} \limsup_{|x| \rightarrow \infty} |x|^{-2} \log(f + 1) < \infty & \quad \text{if } 0 \leq \alpha \leq 1 \\ \limsup_{|x| \rightarrow \infty} |x|^{\frac{2}{\alpha}} f < \infty & \quad \text{if } -1 \leq \alpha < 0 \end{aligned} \quad (5.2)$$

## Theorem

For the interaction rates  $(a + b)^\alpha$  with  $|\alpha| \leq 1$ , suppose  $\{u_i^\varepsilon\}$  are unique mild solutions of (1.1) and have initial data  $g_i \in X_{n,\alpha}$ .

- Then, there exists  $\bar{T} = \bar{T}(n, \alpha, \{g_i\}) \in (0, \infty]$  such that

$$u_i^\varepsilon \rightarrow \frac{\rho}{2n} \text{ in } L^1_{loc}(Q_T) \text{ for all } 0 < T < \bar{T} \text{ as } \varepsilon \rightarrow 0. \quad (5.3)$$

In case  $\alpha = 1$ , if solution of (1.2) have no uniqueness for the initial data  $\sum_{i=1}^{2n} g_i$ , the convergence takes along a subsequence for each given sequence  $\varepsilon_j \rightarrow 0$ . Otherwise, the convergence is arbitrary as  $\varepsilon \rightarrow 0$ .

## continued

- If initial data  $g_i$  has decomposition

$$g_i = l_i + h_i, \quad h_i \in L^1(\mathbb{R}^n), \quad \text{and} \quad l_i \in C^1(\mathbb{R}^n) \quad \text{with} \quad \int |Dl_i| < \infty \quad (5.4)$$

, then

$$\rho^\varepsilon \rightarrow \rho \text{ in } C([0, T], L^1_{loc}(\mathbb{R}^n)) \text{ for all } 0 < T < \bar{T}. \quad (5.5)$$

- When  $\alpha \neq 1$ ,  $\rho \in C([0, T], L^1_{loc}(\mathbb{R}^n))$  is the unique weak solution of  $\partial_t \rho - \nabla \cdot \left( \frac{1}{n^{1-\alpha} \cdot \rho^\alpha} \nabla \rho \right) = 0$  with initial data  $\sum_{i=1}^{2n} g_i$ . In case  $\alpha = 1$ ,  $\rho \in C([0, T], L^1_{loc}(\mathbb{R}^n))$  is some weak solution of  $\partial_t \rho - \nabla \cdot \left( \frac{1}{\rho} \nabla \rho \right) = 0$  with initial data  $\sum_{i=1}^{2n} g_i$ .

- $\bar{T}$  could be taken as  $\bar{T} = C(n, \alpha) \cdot \min(T_1, T_2)$ ,  $T_i \in (0, \infty]$  where








$$T_1 = \begin{cases} [\min_{i=1}^{2n} (\liminf_{|x| \rightarrow \infty} |x|^{\frac{2}{\alpha}} f_i)]^{\frac{1}{\alpha}} & \text{if } \frac{2}{n} \leq \alpha \leq 1 \\ \infty & \text{if } -1 \leq \alpha < \frac{2}{n} \end{cases}$$

and

$$T_2 = \begin{cases} [\max_{i=1}^{2n} \limsup_{|x| \rightarrow \infty} |x|^{-2} \log(f_i + 1)]^{-1} & \text{if } 0 \leq \alpha \leq 1 \\ [\max_{i=1}^{2n} (\limsup_{|x| \rightarrow \infty} |x|^{\frac{2}{\alpha}} f_i)]^{\frac{1}{\alpha}} & \text{if } -1 \leq \alpha < 0. \end{cases}$$

Here, negative power of 0 is defined to be  $\infty$  and  $f_i$  is a function with  $f_i - g_i \in L^1$  given from definition of admissible data.

# References

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# The End