

# Dichotomy theorem for a degenerate one dimensional Keller-Segel model.

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## The classical Keller-Segel model

We are interested in a **population of cells** which attract themselves through a chemical signal.

The PATLAK, KELLER AND SEGEL model involves two species:

- the cell density  $\rho(t, x)$ ,
- the chemoattractant concentration  $c(t, x)$ .

$$\begin{cases} \frac{\partial \rho}{\partial t} = \Delta \rho - \chi \nabla \cdot (\rho \nabla c) & t > 0, x \in \Omega \subset \mathbb{R}^2, \\ -\Delta c = \rho \end{cases}$$

Parameters:

- The **chemosensitivity** coefficient  $\chi$
- The **total mass** of cells  $M$ , which is conserved along time.

In the whole space  $\mathbb{R}^2$ ,

$$c(t, x) = -\frac{1}{2\pi} \int_{\mathbb{R}^2} \log |x - y| \rho(t, y) dy .$$

## Dichotomy theorem in the whole space $\mathbb{R}^2$

### Theorem (Blanchet, Dolbeault and Perthame)

Assume  $\rho_0(|\log \rho_0| + (1 + |x|^2)) \in L^1(\mathbb{R}^2)$ . If  $\chi M < 8\pi$  then solutions are global in time (*dispersion*); they blow up in finite time if  $\chi M > 8\pi$  (*aggregation*).

In case of subcritical regime  $\chi M < 8\pi$ , the solution converges towards a self-similar profile (diffusion-dominating rescaling).

Main tool: **Dissipated free energy**.

$$\searrow \mathcal{F}[\rho] = \int_{\mathbb{R}^2} \rho(y) \log \rho(y) dy + \frac{\chi}{4\pi} \iint_{\mathbb{R}^2 \times \mathbb{R}^2} \rho(x) \log |x-y| \rho(y) dx dy.$$

There are dozens of contributions (prior to this theorem), among which NANJUNDIAH, JÄGER AND LUCKHAUS, NAGAI, BILER, HERRERO AND VELÁZQUEZ, GAJEWSKI AND ZACHARIAS, HORSTMANN, SENBA AND SUZUKI... '92...

## The variance computation

We fix the center of mass equal to zero and introduce the **variance** of the density:

$$\Pi_2(t) = \int |x|^2 \rho(t, x) dx.$$

We impose  $\Pi_2(0) < +\infty$  and compute formally that

$$\begin{aligned} \frac{d}{dt} \Pi_2(t) &= \int |x|^2 \partial_t \rho(t, x) dx = - \int 2x \nabla \rho(t, x) dx \\ &+ \chi \iint 2x \left( \rho(t, x) \frac{-1}{2\pi} \frac{1}{x-y} \rho(t, y) \right) dx dy \\ &= 2d \int \rho(t, x) dx - \frac{\chi}{2\pi} \left( \int \rho(t, x) \right)^2 \\ &= M \left( 4 - \frac{\chi M}{2\pi} \right) \end{aligned}$$

## The degenerate Keller-Segel model

The Keller-Segel model with a **non-linear diffusion**

$$\begin{cases} \partial_t \rho(t, x) = \Delta \rho^m(t, x) - \chi \nabla \cdot (\rho(t, x) \nabla S(t, x)), & t > 0, x \in \mathbb{R}^d, \\ S(t, x) = - \int_{\mathbb{R}^d} |x - y|^{d(1-m)} \rho(t, y) dy. \end{cases}$$

with the associated free energy

$$\begin{aligned} \mathcal{F}_m[\rho] = & \frac{1}{m-1} \int_{\mathbb{R}^d} \rho^m(x) dx \\ & - \chi \frac{1}{d(m-1)} \iint_{\mathbb{R}^d \times \mathbb{R}^d} |x-y|^{d(1-m)} \rho(x) \rho(y) dx dy \end{aligned}$$

Is there a **dichotomy theorem** ?

BLANCHET CARILLO LAURENCOT '09 in dimension  $d \geq 3$ .

## The Variance computation

For  $m > 1$ , with  $\Pi_2(t) = \int |x|^2 \rho(t, x) dx$ ,

$$\frac{d}{dt} \Pi_2(t) = (m-1) \mathcal{F}_m(\rho(t)).$$

$$\frac{d}{dt} \mathcal{F}_m(\rho(t)) \leq 0.$$

- If  $\mathcal{F}_m \geq 0$  there is no contradiction.

Otherwise:

- If  $\mathcal{F}_m(\rho_0) < 0$ , the solution blows-up in finite time.
- If  $\mathcal{F}_m(\rho_0) \geq 0$  ??

The critical parameter is given by

$$\chi_c = \inf [\chi \geq 0 \mid \exists \rho \text{ with } \mathbb{F}_m(\rho) < 0]$$

Do we always have blow-up in the supercritical case ?

# A one dimensional particles scheme

## A gradient flow interpretation

Assume without loss of generality  $M = 1$ .

The Keller-Segel system with a **logarithmic weight**:

$$\partial_t \rho = \Delta \rho + \frac{\chi}{d\pi} \nabla \cdot (\rho \nabla \log |z| * \rho) ,$$

$$\partial_t u = \Delta u + \nabla \cdot (-uy + \frac{\chi}{d\pi} \nabla \log |z| * u) ,$$

is formally the (generalized) **gradient flow** of the free energy functional in the space of probability measures endowed with the **Wasserstein metric**.



## Reformulation in 1D

We reformulate the degenerate Keller-Segel equations in dimension 1.

The energy becomes for any increasing function  $X \in L^2(0, 1)$ :

$$F(X) = \frac{1}{m-1} \int_{(0,1)} \left( \frac{dX}{dp} \right)^{1-m} dp - \frac{\chi}{m-1} \iint_{(0,1)^2} |X(p) - X(q)|^{1-m}, dpdq.$$

## Particles approximation

We take a discretization step  $h_N$  ( $h_N = \frac{1}{N}$ ), and define for  $i = 0..N + 1$ ,  $X_i = X(ih_N)$ . Each particle carries a mass  $h_N$ . It leads to the discrete energy  $F_N : \mathbb{R}^N \rightarrow \mathbb{R}$ :

$$F_N(X) = \sum_{i=1}^{N-1} (X_{i+1} - X_i)^{1-m} - \chi \sum_{1 \leq i \neq j \leq N} |X_i - X_j|^{1-m},$$

and the associated gradient flow equation:

Definition (Gradient flow equation.)

$$\begin{cases} \dot{X}(t) = -\nabla F_N(X(t)) & t \in \mathbb{R} \\ X(0) = X^0 & X^0 \in \mathbb{R}^N. \end{cases}$$

The scheme was studied by DEVYS '10 FOR THE CLASSICAL KELLER-SEGEL CASE.

## Critical parameters

### Definition

Let  $p \in \mathbb{N}$  and

$$C_p = \min \left\{ \sum_{i=1}^{p-1} (X_{i+1} - X_i)^{1-m} \left| \sum_{1 \leq i \neq j \leq p} |X_i - X_j|^{1-m} = 1 \right. \right\}.$$

- $C_p > \frac{1}{p}$  (variant of the Hardy-Littlewood-Sobolev inequality).
- $F_p$  is bounded by below iff  $\chi \leq C_p$ .

# Dichotomy theorem

## Theorem

*Let  $\chi > C_N$ , and  $2 > m > 1$ . Let  $X$  be a solution of the the discrete gradient flow*

- If for all  $p \in [1, N]$ ,  $\chi \neq C_p$  then  $X$  blows up in finite time.*
- If there exists  $p \in [1, N]$  such that  $\chi = C_p$  then either  $X$  blows up in finite or  $Y = \frac{X}{|X|}$  blows up (at infinity) + rigidity.*

## Idea of the proof

We define  $f_{m+1}(X(t)) = |X(t)|^{m+1}$ .

- $f_{m+1}$  is concave.
- $f'_{m+1}(X(t)) = (m+1)(m-1)F_N(X(t))|X(t)|^{m-1}$ .
- $f''_{m+1}(X(t)) = -\frac{(m+1)(m-1)}{f_{m+1}(X(t))} H_{m+1}^N\left(\frac{X(t)}{|X(t)|}\right)$ .

where for any  $Y$  with  $|Y| = 1$ ,

$$H_N(Y) = \left[ |\nabla F_N(Y)|^2 - ((m-1)F_N(Y))^2 \right] \geq 0.$$

- Convergence to a critical point of  $H_N$  ( and  $F_N$  ).
- Critical points of  $F_N$  are minimizers.

Compactness issue.

# Perspectives

- The 1D/2D continuous case.
- Blow-up in finite time.
- Rigidity of the blow-up