

# How far can chemotactic cross-diffusion enforce exceeding carrying capacities?

Michael Winkler  
Universität Paderborn

*Workshop in Nonlinear PDEs*  
Brussels, September 7-11, 2015

# Chemotaxis

Chemotaxis:

- ▶ Cells secrete chemical signal substance
- ▶ Cells partially orient their movement toward increasing signal concentration

# Chemotaxis

Chemotaxis:

- ▶ Cells secrete chemical signal substance
- ▶ Cells partially orient their movement toward increasing signal concentration

Chemotactic movement plays a key role in many processes of communication between cells, e.g. in

- ▶ formation of aggregates such as in populations of *Dictyostelium discoideum* or *Escherichia coli*
- ▶ tumor cell migration
- ▶ organization of cell positioning during embryonic development
- ▶ ...

# The classical Keller-Segel model

KELLER/SEGEL 1970:

- ▶  $u = u(x, t)$ : Density of cell population
- ▶  $v = v(x, t)$ : Concentration of signal

$$u_t = \Delta u - \nabla \cdot (u \nabla v), \quad x \in \Omega, t > 0,$$

$$v_t = \Delta v - v + u, \quad x \in \Omega, t > 0.$$

(KS)

- ▶  $\Omega \subset \mathbb{R}^n$ : bounded domain

# The classical Keller-Segel model

KELLER/SEGEL 1970:

- ▶  $u = u(x, t)$ : Density of cell population
- ▶  $v = v(x, t)$ : Concentration of signal

$$\begin{aligned}u_t &= \Delta u - \nabla \cdot (u \nabla v), & x \in \Omega, t > 0, \\v_t &= \Delta v - v + u, & x \in \Omega, t > 0.\end{aligned}$$

(KS)

- ▶  $\Omega \subset \mathbb{R}^n$ : bounded domain
- ▶ Homogeneous Neumann boundary conditions imply mass conservation:

$$\int_{\Omega} u(x, t) dx = \int_{\Omega} u_0(x) dx$$

# The classical Keller-Segel model

---

## Some results on global well-posedness

- ▶  $n = 1$ : All solutions are global and bounded (OSAKI/YAGI 1997).

# The classical Keller-Segel model

---

## Some results on global well-posedness

- ▶  $n = 1$ : All solutions are global and bounded (OSAKI/YAGI 1997).
- ▶  $n = 2$ : If

$$\|u_0\|_{L^1(\Omega)} < 4\pi,$$

then  $(u, v)$  is global and bounded (NAGAI/SENBA/YOSHIDA 1997).

# The classical Keller-Segel model

## Some results on global well-posedness

- ▶  $n = 1$ : All solutions are global and bounded (OSAKI/YAGI 1997).
- ▶  $n = 2$ : If

$$\|u_0\|_{L^1(\Omega)} < 4\pi,$$

then  $(u, v)$  is global and bounded (NAGAI/SENBA/YOSHIDA 1997).

- ▶  $n \geq 3$ : Let  $\delta > 0$ . Then there exists  $\varepsilon(\delta) > 0$  such that whenever

$$\|u_0\|_{L^{\frac{n}{2+\delta}}(\Omega)} \leq \varepsilon(\delta) \quad \text{and} \quad \|v_0\|_{W^{1,n+\delta}(\Omega)} \leq \varepsilon(\delta),$$

the solution  $(u, v)$  is global and bounded (W. 2010).



# The classical Keller-Segel model

## Some results on global well-posedness

- ▶  $n = 1$ : All solutions are global and bounded (OSAKI/YAGI 1997).
- ▶  $n = 2$ : If

$$\|u_0\|_{L^1(\Omega)} < 4\pi,$$

then  $(u, v)$  is global and bounded (NAGAI/SENBA/YOSHIDA 1997).

- ▶  $n \geq 3$ : Let  $\delta > 0$ . Then there exists  $\varepsilon(\delta) > 0$  such that whenever

$$\|u_0\|_{L^{\frac{n}{2}+\delta}(\Omega)} \leq \varepsilon(\delta) \quad \text{and} \quad \|v_0\|_{W^{1,n+\delta}(\Omega)} \leq \varepsilon(\delta),$$

the solution  $(u, v)$  is global and bounded (W. 2010).

Many boundedness results are available for related systems, involving e.g. nonlinear diffusion or variants in cross-diffusion, signal production,...(NAGAI, SENBA, SUZUKI, YOSHIDA, LAURENÇOT, BILER, FRIEDMAN, PERTHAME, MIMURA, CARRILLO, CALVEZ, CORRIAS, HILLEN, PAINTER, TAO, TELLO, STINNER, SUGIYAMA, KOZONO, ISHIDA, YOKOTA, CIEŚLAK, WRZOSEK, WANG, OSAKI, NAKAGUCHI,...)

# **The classical Keller-Segel model**

---

**Detecting unbounded solutions**

# The classical Keller-Segel model

---

## Detecting unbounded solutions

- ▶ If  $\Omega \subset \mathbb{R}^2$  is a disk, then there exists **at least one** pair  $(u_0, v_0)$  such that  $(u, v)$  blows up in finite time (HERRERO/VELÁZQUEZ 1997).

# The classical Keller-Segel model

---

## Detecting unbounded solutions

- ▶ If  $\Omega \subset \mathbb{R}^2$  is a disk, then there exists **at least one** pair  $(u_0, v_0)$  such that  $(u, v)$  blows up in finite time (HERRERO/VELÁZQUEZ 1997).
- ▶ If  $\Omega \subset \mathbb{R}^2$  is simply connected, then for almost every  $m > 4\pi$  there exist initial data  $(u_0, v_0)$  such that  $\int_{\Omega} u_0 = m$ , and such that  $(u, v)$  blows up **either in finite or infinite time** (HORSTMANN/WANG 2001).

# The classical Keller-Segel model

## Detecting unbounded solutions

- ▶ If  $\Omega \subset \mathbb{R}^2$  is a disk, then there exists **at least one** pair  $(u_0, v_0)$  such that  $(u, v)$  blows up in finite time (HERRERO/VELÁZQUEZ 1997).
- ▶ If  $\Omega \subset \mathbb{R}^2$  is simply connected, then for almost every  $m > 4\pi$  there exist initial data  $(u_0, v_0)$  such that  $\int_{\Omega} u_0 = m$ , and such that  $(u, v)$  blows up **either in finite or infinite time** (HORSTMANN/WANG 2001).
- ▶ If  $n \geq 3$  and  $\Omega \subset \mathbb{R}^n$  is a ball, then **for all  $m > 0$**  one can find  $(u_0, v_0)$  with mass  $\int_{\Omega} u_0 = m$  such that  $(u, v)$  blows up **either in finite or infinite time** (W. 2010).

# The classical Keller-Segel model

## Detecting unbounded solutions

- ▶ If  $\Omega \subset \mathbb{R}^2$  is a disk, then there exists **at least one** pair  $(u_0, v_0)$  such that  $(u, v)$  blows up in finite time (HERRERO/VELÁZQUEZ 1997).
- ▶ If  $\Omega \subset \mathbb{R}^2$  is simply connected, then for almost every  $m > 4\pi$  there exist initial data  $(u_0, v_0)$  such that  $\int_{\Omega} u_0 = m$ , and such that  $(u, v)$  blows up **either in finite or infinite time** (HORSTMANN/WANG 2001).
- ▶ If  $n \geq 3$  and  $\Omega \subset \mathbb{R}^n$  is a ball, then **for all  $m > 0$**  one can find  $(u_0, v_0)$  with mass  $\int_{\Omega} u_0 = m$  such that  $(u, v)$  blows up **either in finite or infinite time** (W. 2010).

Is blow-up *in finite time* a rarely occurring phenomenon?

# The classical Keller-Segel model

---

## Finite-time blow-up

- ▶  $n \geq 3$  Let  $R > 0$  and  $\Omega = B_R(0) \subset \mathbb{R}^n$ . Then for all  $m > 0$  there exist  $(u_0, v_0)$  with mass  $\int_{\Omega} u_0 = m$  such that  $(u, v)$  blows up in finite time (W. 2013).

# The classical Keller-Segel model

## Finite-time blow-up

- ▶  $n \geq 3$  Let  $R > 0$  and  $\Omega = B_R(0) \subset \mathbb{R}^n$ . Then for all  $m > 0$  there exist  $(u_0, v_0)$  with mass  $\int_{\Omega} u_0 = m$  such that  $(u, v)$  blows up in finite time (W. 2013).

Moreover, for any positive radial  $(u_0, v_0)$  one can find smooth positive radial  $u_{0k}$  and  $v_{0k}$  such that  $(u_{0k}, v_{0k}) \rightarrow (u_0, v_0)$  in  $L^p(\Omega) \times W^{1,2}(\Omega)$  for all  $p < \frac{2n}{n+2}$ , but such that the corresponding solutions  $(u_k, v_k)$  blow up in finite time for each  $k$ .



# The classical Keller-Segel model

## Finite-time blow-up

- ▶  $n \geq 3$  Let  $R > 0$  and  $\Omega = B_R(0) \subset \mathbb{R}^n$ . Then for all  $m > 0$  there exist  $(u_0, v_0)$  with mass  $\int_{\Omega} u_0 = m$  such that  $(u, v)$  blows up in finite time (W. 2013).

Moreover, for any positive radial  $(u_0, v_0)$  one can find smooth positive radial  $u_{0k}$  and  $v_{0k}$  such that  $(u_{0k}, v_{0k}) \rightarrow (u_0, v_0)$  in  $L^p(\Omega) \times W^{1,2}(\Omega)$  for all  $p < \frac{2n}{n+2}$ , but such that the corresponding solutions  $(u_k, v_k)$  blow up in finite time for each  $k$ .

- ▶  $n = 2$  Let  $n = 2$ ,  $R > 0$  and  $\Omega = B_R(0) \subset \mathbb{R}^2$ . Then for all  $m > 8\pi$  there exist  $(u_0, v_0)$  with mass  $\int_{\Omega} u_0 = m$  such that  $(u, v)$  blows up in finite time (MIZOGUCHI/W., preprint).

# The classical Keller-Segel model

## Finite-time blow-up

- ▶  $n \geq 3$  Let  $R > 0$  and  $\Omega = B_R(0) \subset \mathbb{R}^n$ . Then for all  $m > 0$  there exist  $(u_0, v_0)$  with mass  $\int_{\Omega} u_0 = m$  such that  $(u, v)$  blows up in finite time (W. 2013).

Moreover, for any positive radial  $(u_0, v_0)$  one can find smooth positive radial  $u_{0k}$  and  $v_{0k}$  such that  $(u_{0k}, v_{0k}) \rightarrow (u_0, v_0)$  in  $L^p(\Omega) \times W^{1,2}(\Omega)$  for all  $p < \frac{2n}{n+2}$ , but such that the corresponding solutions  $(u_k, v_k)$  blow up in finite time for each  $k$ .

- ▶  $n = 2$  Let  $n = 2$ ,  $R > 0$  and  $\Omega = B_R(0) \subset \mathbb{R}^2$ . Then for all  $m > 8\pi$  there exist  $(u_0, v_0)$  with mass  $\int_{\Omega} u_0 = m$  such that  $(u, v)$  blows up in finite time (MIZOGUCHI/W., preprint).

Moreover, for any positive radial  $(u_0, v_0)$  one can find smooth positive radial  $u_{0k}$  and  $v_{0k}$  such that  $(u_{0k}, v_{0k}) \rightarrow (u_0, v_0)$  in  $L^p(\Omega) \times W^{1,2}(\Omega)$  for all  $p \in (0, 1)$ , but such that the corresponding solutions  $(u_k, v_k)$  blow up in finite time for each  $k$ .

# The classical Keller-Segel model

## Finite-time blow-up

- ▶  $n \geq 3$  Let  $R > 0$  and  $\Omega = B_R(0) \subset \mathbb{R}^n$ . Then for all  $m > 0$  there exist  $(u_0, v_0)$  with mass  $\int_{\Omega} u_0 = m$  such that  $(u, v)$  blows up in finite time (W. 2013).

Moreover, for any positive radial  $(u_0, v_0)$  one can find smooth positive radial  $u_{0k}$  and  $v_{0k}$  such that  $(u_{0k}, v_{0k}) \rightarrow (u_0, v_0)$  in  $L^p(\Omega) \times W^{1,2}(\Omega)$  for all  $p < \frac{2n}{n+2}$ , but such that the corresponding solutions  $(u_k, v_k)$  blow up in finite time for each  $k$ .

- ▶  $n = 2$  Let  $n = 2$ ,  $R > 0$  and  $\Omega = B_R(0) \subset \mathbb{R}^2$ . Then for all  $m > 8\pi$  there exist  $(u_0, v_0)$  with mass  $\int_{\Omega} u_0 = m$  such that  $(u, v)$  blows up in finite time (MIZOGUCHI/W., preprint).

Moreover, for any positive radial  $(u_0, v_0)$  one can find smooth positive radial  $u_{0k}$  and  $v_{0k}$  such that  $(u_{0k}, v_{0k}) \rightarrow (u_0, v_0)$  in  $L^p(\Omega) \times W^{1,2}(\Omega)$  for all  $p \in (0, 1)$ , but such that the corresponding solutions  $(u_k, v_k)$  blow up in finite time for each  $k$ .

In particular: All constant steady states  $(u, v) \equiv (m, m)$  are highly unstable.

# The classical Keller-Segel model

## Too simple for reality?

Considerable efforts in modeling attempt to preclude blow-up in chemotaxis models:

# The classical Keller-Segel model

---

## Too simple for reality?

Considerable efforts in modeling attempt to preclude blow-up in chemotaxis models:

- ▶ Inhibited cross-diffusion at large (cell or signal) densities  
(VELÁZQUEZ 2004, OTHMER/STEVENS 1997, HORSTMANN/W. 2005):

$$u_t = \Delta u - \nabla \cdot (S(u, v) \nabla v)$$

# The classical Keller-Segel model

---

## Too simple for reality?

Considerable efforts in modeling attempt to preclude blow-up in chemotaxis models:

- ▶ Inhibited cross-diffusion at large (cell or signal) densities (VELÁZQUEZ 2004, OTHMER/STEVENS 1997, HORSTMANN/W. 2005):

$$u_t = \Delta u - \nabla \cdot (S(u, v) \nabla v)$$

- ▶ Enhanced diffusion at large densities (KOWALCZYK 2005, SENBA/SUZUKI 2006):

$$u_t = \nabla \cdot (D(u) \nabla u) - \nabla \cdot (u \nabla v)$$

# The classical Keller-Segel model

## Too simple for reality?

Considerable efforts in modeling attempt to preclude blow-up in chemotaxis models:

- ▶ Inhibited cross-diffusion at large (cell or signal) densities (VELÁZQUEZ 2004, OTHMER/STEVENS 1997, HORSTMANN/W. 2005):

$$u_t = \Delta u - \nabla \cdot (S(u, v) \nabla v)$$

- ▶ Enhanced diffusion at large densities (KOWALCZYK 2005, SENBA/SUZUKI 2006):

$$u_t = \nabla \cdot (D(u) \nabla u) - \nabla \cdot (u \nabla v)$$

- ▶ Volume-filling effects (PAINTER/HILLEN 2002, TAO/W. 2010, WANG/WRZOSEK/W. 2012):

$$u_t = \nabla \cdot (D(u) \nabla u) - \nabla \cdot (S(u) \nabla v)$$

# The classical Keller-Segel model

## Too simple for reality?

Considerable efforts in modeling attempt to preclude blow-up in chemotaxis models:

- ▶ Inhibited cross-diffusion at large (cell or signal) densities (VELÁZQUEZ 2004, OTHMER/STEVENS 1997, HORSTMANN/W. 2005):

$$u_t = \Delta u - \nabla \cdot (S(u, v) \nabla v)$$

- ▶ Enhanced diffusion at large densities (KOWALCZYK 2005, SENBA/SUZUKI 2006):

$$u_t = \nabla \cdot (D(u) \nabla u) - \nabla \cdot (u \nabla v)$$

- ▶ Volume-filling effects (PAINTER/HILLEN 2002, TAO/W. 2010, WANG/WRZOSEK/W. 2012):

$$u_t = \nabla \cdot (D(u) \nabla u) - \nabla \cdot (S(u) \nabla v)$$

- ▶ Saturation in signal production (MYERSCOUGH/MAINI/PAINTER 1998, OSAKI/NAKAGUCHI 2011):

$$v_t = \Delta v - v + f(u)$$



# The classical Keller-Segel model

## Too simple for reality?

Considerable efforts in modeling attempt to preclude blow-up in chemotaxis models:

- ▶ Inhibited cross-diffusion at large (cell or signal) densities (VELÁZQUEZ 2004, OTHMER/STEVENS 1997, HORSTMANN/W. 2005):

$$u_t = \Delta u - \nabla \cdot (S(u, v) \nabla v)$$

- ▶ Enhanced diffusion at large densities (KOWALCZYK 2005, SENBA/SUZUKI 2006):

$$u_t = \nabla \cdot (D(u) \nabla u) - \nabla \cdot (u \nabla v)$$

- ▶ Volume-filling effects (PAINTER/HILLEN 2002, TAO/W. 2010, WANG/WRZOSEK/W. 2012):

$$u_t = \nabla \cdot (D(u) \nabla u) - \nabla \cdot (S(u) \nabla v)$$

- ▶ Saturation in signal production (MYERSCOUGH/MAINI/PAINTER 1998, OSAKI/NAKAGUCHI 2011):

$$v_t = \Delta v - v + f(u)$$

- ▶ ...

# Chemotaxis with cell kinetics

## Blow-up exclusion by logistic death effects

$$u_t = \Delta u - \nabla \cdot (u \nabla v) + ru - \mu u^2, \quad x \in \Omega, t > 0,$$

$$v_t = \Delta v - v + u, \quad x \in \Omega, t > 0.$$

# Chemotaxis with cell kinetics

## Blow-up exclusion by logistic death effects

$$u_t = \Delta u - \nabla \cdot (u \nabla v) + ru - \mu u^2, \quad x \in \Omega, t > 0,$$

$$v_t = \Delta v - v + u, \quad x \in \Omega, t > 0.$$

- ▶  $n \leq 2$ : For arbitrary  $\mu > 0$ , all solutions are global and bounded (OSAKI/TSUJIKAWA/YAGI/MIMURA 2002, TELLO/W. 2007).

# Chemotaxis with cell kinetics

## Blow-up exclusion by logistic death effects

$$u_t = \Delta u - \nabla \cdot (u \nabla v) + ru - \mu u^2, \quad x \in \Omega, t > 0,$$

$$v_t = \Delta v - v + u, \quad x \in \Omega, t > 0.$$

- ▶  $n \leq 2$ : For arbitrary  $\mu > 0$ , all solutions are global and bounded (OSAKI/TSUJIKAWA/YAGI/MIMURA 2002, TELLO/W. 2007).
- ▶  $n \geq 3$ : If  $\mu > 0$  is sufficiently large, then all solutions are global and bounded (W. 2010, TELLO/W. 2007).

# Chemotaxis with cell kinetics

## Blow-up exclusion by logistic death effects

$$u_t = \Delta u - \nabla \cdot (u \nabla v) + ru - \mu u^2, \quad x \in \Omega, t > 0,$$

$$v_t = \Delta v - v + u, \quad x \in \Omega, t > 0.$$

- ▶  $n \leq 2$ : For arbitrary  $\mu > 0$ , all solutions are global and bounded (OSAKI/TSUJIKAWA/YAGI/MIMURA 2002, TELLO/W. 2007).
- ▶  $n \geq 3$ : If  $\mu > 0$  is sufficiently large, then all solutions are global and bounded (W. 2010, TELLO/W. 2007).
- ▶ The same conclusions hold for the parabolic-elliptic system in which

$$0 = \Delta v - v + u.$$

# Chemotaxis with logistic sources

## Some simulations...

Consider

$$u_t = 0.0148u_{xx} - (uv_x)_x + 0.1u - 0.1u^2, \quad x \in (0,1), t > 0,$$

$$0 = v_{xx} - v + u, \quad x \in (0,1), t > 0,$$

$$u_x(0, t) = u_x(1, t) = v_x(0, t) = v_x(1, t) = 0, \quad t > 0,$$

$$u(x, 0) = 10x^2 \quad x \in (0,1).$$

# Chemotaxis with logistic sources

## Some simulations...

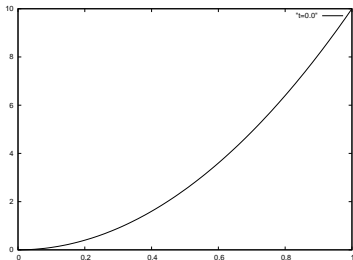
Consider

$$u_t = 0.0148u_{xx} - (uv_x)_x + 0.1u - 0.1u^2, \quad x \in (0,1), t > 0,$$

$$0 = v_{xx} - v + u, \quad x \in (0,1), t > 0,$$

$$u_x(0, t) = u_x(1, t) = v_x(0, t) = v_x(1, t) = 0, \quad t > 0,$$

$$u(x, 0) = 10x^2 \quad x \in (0,1).$$



Solution for  $t = 0$

# Chemotaxis with logistic sources

## Some simulations...

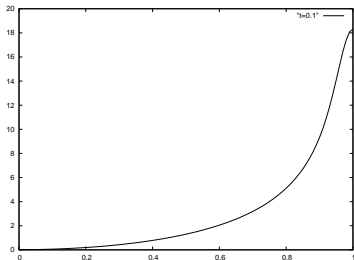
Consider

$$u_t = 0.0148u_{xx} - (uv_x)_x + 0.1u - 0.1u^2, \quad x \in (0,1), t > 0,$$

$$0 = v_{xx} - v + u, \quad x \in (0,1), t > 0,$$

$$u_x(0, t) = u_x(1, t) = v_x(0, t) = v_x(1, t) = 0, \quad t > 0,$$

$$u(x, 0) = 10x^2 \quad x \in (0,1).$$



Solution for  $t = 0.1$



# Chemotaxis with logistic sources

## Some simulations...

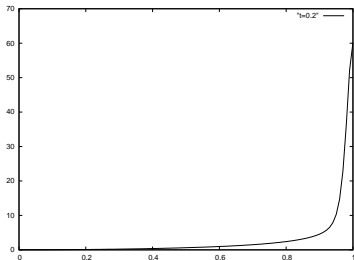
Consider

$$u_t = 0.0148u_{xx} - (uv_x)_x + 0.1u - 0.1u^2, \quad x \in (0,1), t > 0,$$

$$0 = v_{xx} - v + u, \quad x \in (0,1), t > 0,$$

$$u_x(0, t) = u_x(1, t) = v_x(0, t) = v_x(1, t) = 0, \quad t > 0,$$

$$u(x, 0) = 10x^2 \quad x \in (0,1).$$



Solution for  $t = 0.2$

# Chemotaxis with logistic sources

## Some simulations...

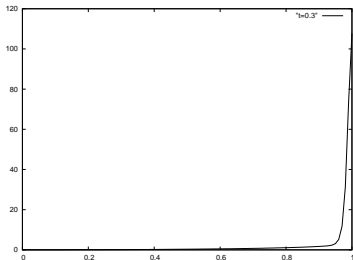
Consider

$$u_t = 0.0148u_{xx} - (uv_x)_x + 0.1u - 0.1u^2, \quad x \in (0,1), t > 0,$$

$$0 = v_{xx} - v + u, \quad x \in (0,1), t > 0,$$

$$u_x(0, t) = u_x(1, t) = v_x(0, t) = v_x(1, t) = 0, \quad t > 0,$$

$$u(x, 0) = 10x^2 \quad x \in (0,1).$$



Solution for  $t = 0.3$

# Chemotaxis with logistic sources

## Some simulations...

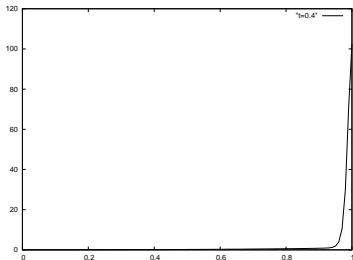
Consider

$$u_t = 0.0148u_{xx} - (uv_x)_x + 0.1u - 0.1u^2, \quad x \in (0,1), t > 0,$$

$$0 = v_{xx} - v + u, \quad x \in (0,1), t > 0,$$

$$u_x(0, t) = u_x(1, t) = v_x(0, t) = v_x(1, t) = 0, \quad t > 0,$$

$$u(x, 0) = 10x^2 \quad x \in (0,1).$$



Solution for  $t = 0.4$

# Chemotaxis with logistic sources

## Some simulations...

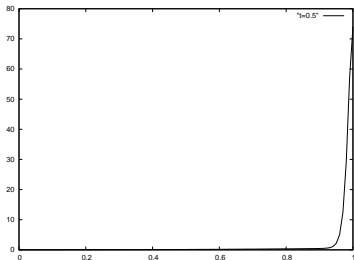
Consider

$$u_t = 0.0148u_{xx} - (uv_x)_x + 0.1u - 0.1u^2, \quad x \in (0,1), t > 0,$$

$$0 = v_{xx} - v + u, \quad x \in (0,1), t > 0,$$

$$u_x(0, t) = u_x(1, t) = v_x(0, t) = v_x(1, t) = 0, \quad t > 0,$$

$$u(x, 0) = 10x^2 \quad x \in (0,1).$$



Solution for  $t = 0.5$

# Chemotaxis with logistic sources

## Some simulations...

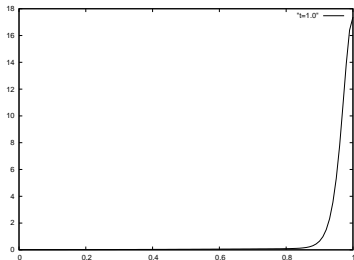
Consider

$$u_t = 0.0148u_{xx} - (uv_x)_x + 0.1u - 0.1u^2, \quad x \in (0,1), t > 0,$$

$$0 = v_{xx} - v + u, \quad x \in (0,1), t > 0,$$

$$u_x(0, t) = u_x(1, t) = v_x(0, t) = v_x(1, t) = 0, \quad t > 0,$$

$$u(x, 0) = 10x^2 \quad x \in (0,1).$$



Solution for  $t = 1.0$

# Chemotaxis with logistic sources

## Some simulations...

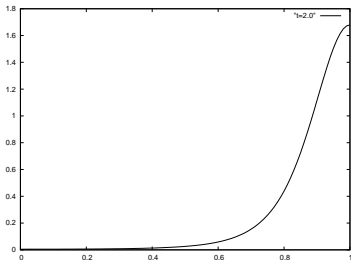
Consider

$$u_t = 0.0148u_{xx} - (uv_x)_x + 0.1u - 0.1u^2, \quad x \in (0,1), t > 0,$$

$$0 = v_{xx} - v + u, \quad x \in (0,1), t > 0,$$

$$u_x(0, t) = u_x(1, t) = v_x(0, t) = v_x(1, t) = 0, \quad t > 0,$$

$$u(x, 0) = 10x^2 \quad x \in (0,1).$$



Solution for  $t = 2.0$

# Chemotaxis with logistic sources

## Some simulations...

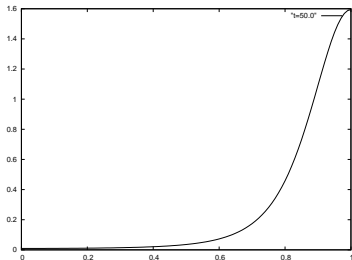
Consider

$$u_t = 0.0148u_{xx} - (uv_x)_x + 0.1u - 0.1u^2, \quad x \in (0,1), t > 0,$$

$$0 = v_{xx} - v + u, \quad x \in (0,1), t > 0,$$

$$u_x(0, t) = u_x(1, t) = v_x(0, t) = v_x(1, t) = 0, \quad t > 0,$$

$$u(x, 0) = 10x^2 \quad x \in (0,1).$$



Solution for  $t = 50.0$

# Chemotaxis with logistic sources

...and a new challenge: A **transient** growth phenomenon

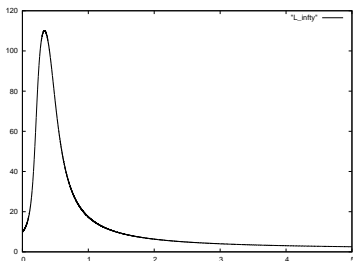
Consider

$$u_t = 0.0148u_{xx} - (uv_x)_x + 0.1u - 0.1u^2, \quad x \in (0,1), t > 0,$$

$$0 = v_{xx} - v + u, \quad x \in (0,1), t > 0,$$

$$u_x(0, t) = u_x(1, t) = v_x(0, t) = v_x(1, t) = 0, \quad t > 0,$$

$$u(x, 0) = 10x^2 \quad x \in (0,1).$$



Evolution of  $\|u(\cdot, t)\|_{L^\infty(\Omega)}$  for  $0 \leq t \leq 5.0$



# Going beyond carrying capacities

## A refined view upon cross-diffusive destabilization

Recall: For

$$u_t = ru - \mu u^2 \quad (1)$$

there exists a **carrying capacity**  $u_c := \frac{r}{\mu}$  such that  $u \rightarrow u_c$  as  $t \rightarrow \infty$   
and

$$u \leq \max\{u_0, u_c\} \quad \text{for all } (x \text{ and}) t.$$

# Going beyond carrying capacities

## A refined view upon cross-diffusive destabilization

Recall: For

$$u_t = ru - \mu u^2 \quad (1)$$

there exists a **carrying capacity**  $u_c := \frac{r}{\mu}$  such that  $u \rightarrow u_c$  as  $t \rightarrow \infty$   
and

$$u \leq \max\{u_0, u_c\} \quad \text{for all } (x \text{ and}) t.$$

The same holds for

$$u_t = \Delta u + ru - \mu u^2. \quad (2)$$

# Going beyond carrying capacities

## A refined view upon cross-diffusive destabilization

Recall: For

$$u_t = ru - \mu u^2 \quad (1)$$

there exists a **carrying capacity**  $u_c := \frac{r}{\mu}$  such that  $u \rightarrow u_c$  as  $t \rightarrow \infty$   
and

$$u \leq \max\{u_0, u_c\} \quad \text{for all } (x \text{ and}) t.$$

The same holds for

$$u_t = \Delta u + ru - \mu u^2. \quad (2)$$

That is: In both (1) and (2):

Populations cannot exceed the carrying capacity *dynamically*.

# Going beyond carrying capacities

## A refined view upon cross-diffusive destabilization

Consider

$$u_t = \varepsilon u_{xx} - (uv_x)_x + ru - \mu u^2, \quad x \in \Omega, t > 0,$$

$$0 = v_{xx} - v + u, \quad x \in \Omega, t > 0.$$

in  $\Omega = (0, L) \subset \mathbb{R}$ , with  $\varepsilon > 0$ ,  $r \geq 0$  and  $\mu > 0$ .

# Going beyond carrying capacities

## A refined view upon cross-diffusive destabilization

Consider

$$\begin{aligned} u_t &= \varepsilon u_{xx} - (uv_x)_x + ru - \mu u^2, & x \in \Omega, t > 0, \\ 0 &= v_{xx} - v + u, & x \in \Omega, t > 0. \end{aligned}$$

in  $\Omega = (0, L) \subset \mathbb{R}$ , with  $\varepsilon > 0$ ,  $r \geq 0$  and  $\mu > 0$ .

All solutions are global and bounded (TELLO/W. 2007).

# Going beyond carrying capacities

## A refined view upon cross-diffusive destabilization

Consider

$$\begin{aligned} u_t &= \varepsilon u_{xx} - (uv_x)_x + ru - \mu u^2, & x \in \Omega, t > 0, \\ 0 &= v_{xx} - v + u, & x \in \Omega, t > 0. \end{aligned}$$

in  $\Omega = (0, L) \subset \mathbb{R}$ , with  $\varepsilon > 0$ ,  $r \geq 0$  and  $\mu > 0$ .

All solutions are global and bounded (TELLO/W. 2007).

How far can solutions exceed the carrying capacity  $u_c$  dynamically?

# Going beyond carrying capacities

## A refined view upon cross-diffusive destabilization

Consider

$$u_t = \varepsilon u_{xx} - (uv_x)_x + ru - \mu u^2, \quad x \in \Omega, t > 0,$$

$$0 = v_{xx} - v + u, \quad x \in \Omega, t > 0.$$

in  $\Omega = (0, L) \subset \mathbb{R}$ , with  $\varepsilon > 0$ ,  $r \geq 0$  and  $\mu > 0$ .

# Going beyond carrying capacities

## A refined view upon cross-diffusive destabilization

Consider

$$\begin{aligned} u_t &= \varepsilon u_{xx} - (uv_x)_x + ru - \mu u^2, & x \in \Omega, t > 0, \\ 0 &= v_{xx} - v + u, & x \in \Omega, t > 0. \end{aligned}$$

in  $\Omega = (0, L) \subset \mathbb{R}$ , with  $\varepsilon > 0$ ,  $r \geq 0$  and  $\mu > 0$ .

**Theorem (W.):** Let  $\mu < 1$ . Then for all  $p > \frac{1}{1-\mu}$  there exists  $C(p) > 0$  such that whenever

$$\|u_0\|_{L^p(\Omega)} > C(p) \cdot \max \left\{ \frac{1}{|\Omega|} \int_{\Omega} u_0, \frac{r}{\mu} \right\},$$

one can find  $T > 0$  such that for all  $M > 0$  there exists  $\varepsilon_0(M) > 0$  with the following property:

For any  $\varepsilon \in (0, \varepsilon_0(M))$  one can pick  $t_\varepsilon \in (0, T)$  and  $x_\varepsilon \in \Omega$  such that

$$u(x_\varepsilon, t_\varepsilon) > M.$$



# Outline of proof

## Analyzing a hyperbolic-elliptic limit system

Consider the case  $\varepsilon = 0$ , that is,

$$u_t = -(uv_x)_x + ru - \mu u^2, \quad x \in \Omega, t > 0,$$

$$0 = v_{xx} - v + u, \quad x \in \Omega, t > 0.$$

# Outline of proof

## Analyzing a hyperbolic-elliptic limit system

Consider the case  $\varepsilon = 0$ , that is,

$$\begin{aligned} u_t &= -(uv_x)_x + ru - \mu u^2, & x \in \Omega, t > 0, \\ 0 &= v_{xx} - v + u, & x \in \Omega, t > 0. \end{aligned}$$

**Lemma 1:** Let  $q > 1$  and  $u_0 \in W^{1,q}(\Omega)$ . Then there exist  $T_{\max} \leq \infty$  and a generalized solution for  $t \in (0, T_{\max})$  such that for all  $q > 1$ ,

$$u \in L^\infty((0, T); W^{1,q}(\Omega)) \quad \text{for all } T < T_{\max},$$

and such that

$$\text{if } T_{\max} < \infty \quad \text{then} \quad \limsup_{t \nearrow T_{\max}} \|u(\cdot, t)\|_{L^\infty(\Omega)} = \infty.$$

# Outline of proof

## Analyzing a hyperbolic-elliptic limit system

Consider the case  $\varepsilon = 0$ , that is,

$$\begin{aligned} u_t &= -(uv_x)_x + ru - \mu u^2, & x \in \Omega, t > 0, \\ 0 &= v_{xx} - v + u, & x \in \Omega, t > 0. \end{aligned}$$

**Lemma 1:** Let  $q > 1$  and  $u_0 \in W^{1,q}(\Omega)$ . Then there exist  $T_{\max} \leq \infty$  and a generalized solution for  $t \in (0, T_{\max})$  such that for all  $q > 1$ ,

$$u \in L^\infty((0, T); W^{1,q}(\Omega)) \quad \text{for all } T < T_{\max},$$

and such that

$$\text{if } T_{\max} < \infty \quad \text{then} \quad \limsup_{t \nearrow T_{\max}} \|u(\cdot, t)\|_{L^\infty(\Omega)} = \infty.$$

Proof: Construct  $(u, v)$  as a limit of solutions  $(u_\varepsilon, v_\varepsilon)$  of the above problem.

# Outline of proof

## Analyzing a hyperbolic-elliptic limit system

Consider the case  $\varepsilon = 0$ , that is,

$$u_t = -(uv_x)_x + ru - \mu u^2, \quad x \in \Omega, t > 0,$$

$$0 = v_{xx} - v + u, \quad x \in \Omega, t > 0.$$

**Lemma 1:** Let  $q > 1$  and  $u_0 \in W^{1,q}(\Omega)$ . Then there exist  $T_{\max} \leq \infty$  and a generalized solution for  $t \in (0, T_{\max})$  such that for all  $q > 1$ ,

$$u \in L^\infty((0, T); W^{1,q}(\Omega)) \quad \text{for all } T < T_{\max},$$

and such that

$$\text{if } T_{\max} < \infty \quad \text{then} \quad \limsup_{t \nearrow T_{\max}} \|u(\cdot, t)\|_{L^\infty(\Omega)} = \infty.$$

Proof: Construct  $(u, v)$  as a limit of solutions  $(u_\varepsilon, v_\varepsilon)$  of the above problem.

Main Estimate:

$$\int_{\Omega} |u_x|^q \leq \left\{ \int_{\Omega} |u_{0x}|^q + q \int_0^t \left( \int_{\Omega} u(\cdot, s) \right)^{q+1} ds \right\} \cdot \exp \left( 4q \int_0^t \|u(\cdot, s)\|_{L^\infty(\Omega)} ds + qrt \right)$$

# Outline of proof

## Analyzing a hyperbolic-elliptic limit system

Formally, for any  $p > 1$ ,

$$\frac{1}{p} \frac{d}{dt} \int_{\Omega} u^p = - \int_{\Omega} (uv_x)_x u^{p-1} + r \int_{\Omega} u^p - \mu \int_{\Omega} u^{p+1}$$

# Outline of proof

## Analyzing a hyperbolic-elliptic limit system

Formally, for any  $p > 1$ ,

$$\begin{aligned}\frac{1}{p} \frac{d}{dt} \int_{\Omega} u^p &= - \int_{\Omega} (uv_x)_x u^{p-1} + r \int_{\Omega} u^p - \mu \int_{\Omega} u^{p+1} \\ &= (p-1) \int_{\Omega} u^{p-1} u_x v_x + r \int_{\Omega} u^p - \mu \int_{\Omega} u^{p+1}\end{aligned}$$

# Outline of proof

## Analyzing a hyperbolic-elliptic limit system

Formally, for any  $p > 1$ ,

$$\begin{aligned}\frac{1}{p} \frac{d}{dt} \int_{\Omega} u^p &= - \int_{\Omega} (uv_x)_x u^{p-1} + r \int_{\Omega} u^p - \mu \int_{\Omega} u^{p+1} \\ &= (p-1) \int_{\Omega} u^{p-1} u_x v_x + r \int_{\Omega} u^p - \mu \int_{\Omega} u^{p+1} \\ &= -\frac{p-1}{p} \int_{\Omega} u^p v_{xx} + r \int_{\Omega} u^p - \mu \int_{\Omega} u^{p+1}\end{aligned}$$

# Outline of proof

## Analyzing a hyperbolic-elliptic limit system

Formally, for any  $p > 1$ ,

$$\begin{aligned}\frac{1}{p} \frac{d}{dt} \int_{\Omega} u^p &= - \int_{\Omega} (uv_x)_x u^{p-1} + r \int_{\Omega} u^p - \mu \int_{\Omega} u^{p+1} \\ &= (p-1) \int_{\Omega} u^{p-1} u_x v_x + r \int_{\Omega} u^p - \mu \int_{\Omega} u^{p+1} \\ &= -\frac{p-1}{p} \int_{\Omega} u^p v_{xx} + r \int_{\Omega} u^p - \mu \int_{\Omega} u^{p+1} \\ &= -\frac{p-1}{p} \int_{\Omega} u^p v + \frac{p-1}{p} \int_{\Omega} u^{p+1} + r \int_{\Omega} u^p - \mu \int_{\Omega} u^{p+1},\end{aligned}$$

that is,

$$\frac{d}{dt} \int_{\Omega} u^p = pr \int_{\Omega} u^p - (1-p+\mu p) \int_{\Omega} u^{p+1} - (p-1) \int_{\Omega} u^p v.$$



# Outline of proof

## Analyzing a hyperbolic-elliptic limit system

Formally, for any  $p > 1$ ,

$$\begin{aligned}\frac{1}{p} \frac{d}{dt} \int_{\Omega} u^p &= - \int_{\Omega} (uv_x)_x u^{p-1} + r \int_{\Omega} u^p - \mu \int_{\Omega} u^{p+1} \\ &= (p-1) \int_{\Omega} u^{p-1} u_x v_x + r \int_{\Omega} u^p - \mu \int_{\Omega} u^{p+1} \\ &= -\frac{p-1}{p} \int_{\Omega} u^p v_{xx} + r \int_{\Omega} u^p - \mu \int_{\Omega} u^{p+1} \\ &= -\frac{p-1}{p} \int_{\Omega} u^p v + \frac{p-1}{p} \int_{\Omega} u^{p+1} + r \int_{\Omega} u^p - \mu \int_{\Omega} u^{p+1},\end{aligned}$$

that is,

$$\frac{d}{dt} \int_{\Omega} u^p = pr \int_{\Omega} u^p - (1-p+\mu p) \int_{\Omega} u^{p+1} - (p-1) \int_{\Omega} u^p v.$$

Since the rightmost term is nonpositive:

**Corollary 2:** If  $\mu \geq 1$ , then all solutions are global and bounded.

# Outline of proof

## Analyzing a hyperbolic-elliptic limit system

Procedure for  $\mu < 1$ :

- ▶ Derive rigorous justification for

$$\frac{d}{dt} \int_{\Omega} u^p = pr \int_{\Omega} u^p + (p - \mu p - 1) \int_{\Omega} u^{p+1} - (p - 1) \int_{\Omega} u^p v.$$

- ▶ Absorb last term appropriately.

# Outline of proof

## Analyzing a hyperbolic-elliptic limit system

Procedure for  $\mu < 1$ :

- ▶ Derive rigorous justification for

$$\frac{d}{dt} \int_{\Omega} u^p = pr \int_{\Omega} u^p + (p - \mu p - 1) \int_{\Omega} u^{p+1} - (p - 1) \int_{\Omega} u^p v.$$

- ▶ Absorb last term appropriately.

Result:

$$\int_{\Omega} u^p(\cdot, t) \geq C_1 \int_0^t \int_{\Omega} u^{p+1} + \int_{\Omega} u_0^p - C_2 \int_0^t \left( \int_{\Omega} u \right)^{p+1}$$

# Outline of proof

## Analyzing a hyperbolic-elliptic limit system

$$\int_{\Omega} u^p(\cdot, t) \geq C_1 \int_0^t \int_{\Omega} u^{p+1} + \int_{\Omega} u_0^p - C_2 \int_0^t \left( \int_{\Omega} u \right)^{p+1}.$$

# Outline of proof

## Analyzing a hyperbolic-elliptic limit system

$$\int_{\Omega} u^p(\cdot, t) \geq C_1 \int_0^t \int_{\Omega} u^{p+1} + \int_{\Omega} u_0^p - C_2 \int_0^t \left( \int_{\Omega} u \right)^{p+1}.$$

Since  $\int_{\Omega} u^{p+1} \geq C(\int_{\Omega} u^p)^{\frac{p+1}{p}}$ , by a Grönwall-type lemma this implies a blow-up criterion:

**Lemma 3.** Let  $\mu \in (0, 1)$  and  $(1 - \mu)p > 1$ . Then there exists  $C(p) > 0$  with the following property: Whenever  $q > 1$  and  $u_0 \in W^{1,q}(\Omega)$  is nonnegative and such that

$$\|u_0\|_{L^p(\Omega)} > C(p) \cdot \max \left\{ \frac{1}{|\Omega|} \int_{\Omega} u_0, \frac{r}{\mu} \right\},$$

the solution  $(u, v)$  of the hyperbolic-elliptic problem blows up in finite time; that is, we have  $T_{max} < \infty$  and

$$\limsup_{t \nearrow T_{max}} \|u(\cdot, t)\|_{L^\infty(\Omega)} = \infty.$$

# Outline of proof

## Analyzing a hyperbolic-elliptic limit system

By

- ▶ **uniqueness** in the hyperbolic-elliptic problem and – thereby implied –
- ▶ continuous dependence of solutions on the parameter  $\varepsilon \in [0, 1)$ :

# Outline of proof

## Analyzing a hyperbolic-elliptic limit system

By

- ▶ **uniqueness** in the hyperbolic-elliptic problem and – thereby implied –
- ▶ continuous dependence of solutions on the parameter  $\varepsilon \in [0, 1)$ :

The corresponding solutions  $(u_\varepsilon, v_\varepsilon)$  of the parabolic-elliptic problem satisfy

$$u_\varepsilon \rightarrow u \quad \text{as } \varepsilon \searrow 0.$$

# Outline of proof

## Analyzing a hyperbolic-elliptic limit system

By

- ▶ **uniqueness** in the hyperbolic-elliptic problem and – thereby implied –
- ▶ continuous dependence of solutions on the parameter  $\varepsilon \in [0, 1)$ :

The corresponding solutions  $(u_\varepsilon, v_\varepsilon)$  of the parabolic-elliptic problem satisfy

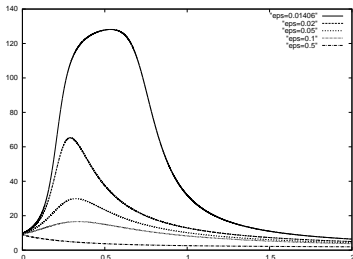
$$u_\varepsilon \rightarrow u \quad \text{as } \varepsilon \searrow 0.$$

Thus, if  $u(x_0, t_0) > 2M$  then  $u_\varepsilon(x_0, t_0) > M$  for all small  $\varepsilon$ .



# Exceeding carrying capacities

A transient small diffusion phenomenon



Solution for  $r = 0.1$ ,  $\mu = 0.1$ ,  $u_0(x) = 10x^2$

Evolution of  $\|u(\cdot, t)\|_{L^\infty(\Omega)}$  for  $0 \leq t \leq 2.0$

Comparing different values of  $\varepsilon$ :  $\varepsilon \in \{0.01406, 0.02, 0.05, 0.1, 0.5\}$

# Conclusion

Chemotactic cross-diffusion can enforce dynamical exceedance of carrying capacities

- ▶ **to an arbitrary extent**, provided that diffusion is sufficiently weak;
- ▶ **even in spatially one-dimensional systems**, where solutions always are bounded.

# Conclusion

Chemotactic cross-diffusion can enforce dynamical exceedance of carrying capacities

- ▶ **to an arbitrary extent**, provided that diffusion is sufficiently weak;
- ▶ **even in spatially one-dimensional systems**, where solutions always are bounded.

Open problems concern e.g.

- ▶ **higher-dimensional** analogues: Does there exist an adequate solution theory for the hyperbolic-elliptic limit? (LANKEIT 2014)
- ▶ **blow-up** in (three-dimensional) system with diffusion and small logistic absorption;
- ▶ problems involving **degenerate diffusion** instead of a constantly small diffusivity.

# Conclusion

Chemotactic cross-diffusion can enforce dynamical exceedance of carrying capacities

- ▶ **to an arbitrary extent**, provided that diffusion is sufficiently weak;
- ▶ **even in spatially one-dimensional systems**, where solutions always are bounded.

Open problems concern e.g.

- ▶ **higher-dimensional** analogues: Does there exist an adequate solution theory for the hyperbolic-elliptic limit? (LANKEIT 2014)
- ▶ **blow-up** in (three-dimensional) system with diffusion and small logistic absorption;
- ▶ problems involving **degenerate diffusion** instead of a constantly small diffusivity.

**Thank you!**