

Orbital Stability for the Schrödinger Operator Involving Inverse Square Potential

G. P. Trachanas and N. B. Zographopoulos

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$$-\Delta u - \left(\frac{N-2}{2}\right)^2 \frac{u}{|x|^2} + \lambda u - |u|^{q-2}u = 0, \quad x \in \mathbb{R}^N. \quad (2)$$

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- Interpolation inequalities ?

Functional Space - Properties

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The functional space H is defined with the norm

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$$\Lambda_\varepsilon(u) := \frac{N-2}{2} \varepsilon^{-1} \int_{|x|=\varepsilon} |u|^2 dS,$$

represents the *Hardy (hidden) energy at the singularity*.

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Lemma

Let H_r be the subspace of H consisting of radial functions. Then, the inclusion

$$H_r \hookrightarrow L^q(\mathbb{R}^N), \quad 2 < q < 2^*,$$

is compact.

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As an interpolation inequality, directly **use the CKN inequality** ($\alpha = \beta = \gamma = -\frac{N-2}{2}$), which provides the same upper bound for q as in the classical case. ■

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Λ may be [the main part of the energy](#).

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where g decays to zero at infinity.

References

References

-  G. P. Trachanas, N. B. Z, *Orbital stability for the Schrödinger operator involving inverse square potential*, J. Diff. Eq., 259, (2015), 4989-5016.
-  J. L. Vázquez, N. B. Z, *Functional aspects of the Hardy inequality: appearance of a hidden energy*, J. Evol. Equ. **12** (2012), 713-739.
-  J. L. Vázquez, N. B. Z, *Hardy type inequalities and hidden energies*, Discrete Contin. Dyn. Syst. 33 (2013), 5457-5491.
-  N. B. Z, *Existence of extremal functions for a Hardy-Sobolev inequality*, J. Funct. Anal. **259** (2010), 308-314.
-  N. B. Z, *Some Results Concerning Hardy and Hardy Type Inequalities*, Handbook of Functional Equations, Springer Optimization and Its Applications, Volume 95, 2014, 535-555.